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Saulnier, Ivan Panaiotov

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#### ORIGINAL CONTRIBUTION

# Interfacial behavior of lipid nanocapsules spread on model membrane monolayers

Tz Ivanova • K. Mircheva • K. Balashev • I. Minkov • P. Saulnier • I. Panaiotov

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Abstract The lipid nanocapsules (LNCs) spread at the airwater interface (A/W) undergo destabilization and disaggregation leading to formation of a triglyceride (TG) surface film. The kinetics of reorganization and formation of TG surface film were followed by measuring either the change of surface pressure at constant area or the surface area at constant surface pressure. From the obtained experimental data were determined the effectiveness of TG spreading and the rate of LNC disaggregation at A/W interface covered with preformed model membrane monolayers of DPPC, Curosurf®, and mucus. Partial LNC stabilization due to their interaction with the model membrane monolayers was observed and characterized by atomic force microscopy (AFM). The obtained results demonstrated that the LNCs spread on mucus surface layer, which models the epithelial surface were more stable than if they were spread either on DPPC or Curosurf<sup>®</sup> surface layers, which emulate the alveolar surface.

**Keywords** Lipid nanocapsules · Spreading · Model membrane monolayers · Atomic force microscopy (AFM)

#### Abbreviations

LNCs	Lipid nanocapsules
TG	Triglyceride
Lab≡TG	Labrafac®
DPPC	Dipalmitoylphosphatidylcholine
A/W	Air-water
LS	Lung surfactant

T. Ivanova (🖂) · K. Mircheva · K. Balashev · I. Minkov · I. Panaiotov

Biophysical Chemistry Laboratory, Dept. of Physical Chemistry, Faculty of Chemistry and Pharmacy, University of Sofia, J. Bourchier str. 1, 1164 Sofia, Bulgaria e-mail: tzivanova@chem.uni-sofia.bg

#### P. Saulnier

INSERM U1066 MINT "Micro et Nanomédecines Biomimétiques, LUNAM Université, 49100 Angers, France

#### Introduction

Over the last decade after the introduction of basic or modified lipid nanocapsules (LNCs) in pharmaceutical practice, they have been successfully utilized as effective drug delivery systems [1-6]. The basic LNCs obtained by phase-inversion method contain a triglyceride (TG) core covered by phospholipid molecules and a soft layer formed from hydroxystearate of polyethylene glycol [1, 2]. The mechanisms of destabilization and reorganization of LNCs after their spreading at pure air-water (A/W) interface were studied in various aspects whereby two populations of nanocapsules without (LNCs I) and with (LNCs II) phospholipid molecules were identified [7–10]. Both fractions were found to be stable and to remain intact in the bulk of their suspensions. A loss of mechanical stability at A/W interface was observed for the fraction without lipid molecules, which underwent rapid destabilization with a rate constant  $k_1$  followed by disaggregation and formation of a TG surface film. It was assumed that  $k_1$  was related to the characteristic time of interfacial disaggregation of the unstable type I nanoparticles, whereas the spreading of TG molecules confined into capsules was instantaneous.

Studying the mechanisms of loss of mechanical stability and reorganization of LNCs containing encapsulated liposoluble drugs in the processes of interaction with various model membrane systems seems appropriate for better understanding of their behavior at the cell membrane interface. Thus, using a monolayer spread at the A/W interface as a simplified convenient membrane model, three monolayer systems were chosen : (*i*) Dipalmitoylphosphatidylcholine (DPPC) monolayer mimicking the cell membrane; (*ii*) lung surfactant (LS) layer accumulated at the alveolar surface of lung; and (*iii*) pig intestinal mucus layer.

The monolayer of saturated phospholipid DPPC is largely preferred as a model of biomembrane because DPPC is the most common lipid component of cell membranes [11–15] as well as because the drug insertion into preformed DPPC monolayers is found to be alike as drug penetration via

biological membranes [16]. DPPC is also the main component of LS. The physiological contribution of LS accumulated at the alveolar surface is related to the maintenance of structural stability of the alveoli during respiration process by reducing the surface tension between alveolar lining layer and gas phase. Curosurf<sup>®</sup> is a commercially available natural LS, which contains about 70 % DPPC of total phospholipid content. It is clinically applicable in cases of LS absence in premature infants [17, 18]. Both DPPC and Curosurf<sup>®</sup> monolayers were often used as model systems of the alveolar surface. For example, the insertion of pulmonary antitubercular drug into LS monolayer has been investigated as the drug was administrated by an appropriated drug delivery system toward the model alveolar surface [19].

The epithelial surface is covered by a highly glycosylated mucus layer that serves as a physical barrier between the extracellular environment and plasma membranes [20]. Recently, the interaction of newly developed LNCs for oral delivery with pig intestinal mucus layer having thickness of several hundred microns were studied by fluorescence resonance energy transfer (FRET) method in order to improve the bioavailability of liposoluble drug Paclitaxel encapsulated in these LNCs [21]. Investigating the processes of adsorption on mucus layer also requires finding a model system where the experimental conditions and physicochemical parameters are under strict control [22, 23]. As such, the mucus monolayer at A/W interface appears to be the right candidate.

The purpose of this paper is to study the mechanisms of destabilization and reorganization of basic LNCs spread on three types preformed at the A/W interface monolayers of: DPPC, Curosurf<sup>®</sup> and pig intestinal mucus. In this approach, the monolayers of DPPC and Curosurf<sup>®</sup> served as a model of the alveolar surface and the mucus monolayer is considered as a model of the luminal surface of gastrointestinal tract. Furthermore, the monolayers studies at A/W interface were accompanied with atomic force microscopy (AFM) imaging of Langmuir–Blodgett (LB) films of the monolayer systems under investigation after their transfer from the A/W interface on mica solid supports.

#### Materials and methods

#### Materials

DPPC with a molecular weight (MW) of 734.04 was provided by Sigma-Aldrich. Labrafac<sup>®</sup> WL1349 (Lab) provided by Gattefossé S.A. (Saint- Priest, France) is a mixture of caprylic and capric acid TG with average MW=512. Lipoid<sup>®</sup> (Lip) S75-3 from Lipoid GmbH (Ludwigshafen, Germany) is a soybean lecithin with 69 % of phosphatidylcholine and an average MW=800. Solutol<sup>®</sup> HS15 (Sol) provided by BASF (Ludwigshafen, Germany) with an average MW=911 is a mixture of free polyethylene glycol 660 (~30 %) and 12-hydroxystearate of polyethylene glycol 660 (~70 %). Curosurf<sup>®</sup> (Chiesi Farmaceutici, Parma, Italy) is a natural surfactant extract, prepared from porcine lungs, containing almost exclusively polar lipids in particular phosphatidylcholine (about 70 % of total phospholipid content), phosphatidylglycerol (about 30 % of total phospholipid content) and about 1 % of surfactant associated hydrophobic proteins SP-B and SP-C. Pig intestinal mucus containing mucin 5 %, lipids 37 %, proteins 39 %, DNA 6 % other 13 % of dry weight was prepared by the colleagues from the University of Angers (France). NaCl was a product of Theokom (Sofia, Bulgaria). Double distilled water was used in all experiments.

Considering the complex composition of each product, the brand names and their abbreviations will be used in the article.

#### Methods

#### Preparation of LNC

LNCs were prepared using the procedure described by Heurtault et al. [1, 2]. All components (Labrafac<sup>®</sup>, Lipoid<sup>®</sup>, Solutol<sup>®</sup>, NaCl, and doubly distilled water) were mixed. Three cycles of progressive heating and cooling combined with magnetic stirring at an agitation speed of 200 rpm, were performed. The first and second cycles were between 85 and 60 °C, whereas the third one was between 85 and 70 °C. When the sample temperature reached 70 °C, a fast cooling dilution was made by rapid addition of 12.5 cm<sup>3</sup> cold ( $\sim T=0$  °C) double-distilled water. After preparation of the LNC aqueous suspension, it was dialyzed against pure water for 48 h in order to take away the free molecules. Spectra/Por<sup>®</sup> Biotech Cellulose Ester Dialysis Membrane (MW=10,000 Da) was used.

The size of LNCs was determined by photon correlation spectroscopy method using a Malvern autosizer<sup>®</sup> 4700 (Malvern Instruments Ltd., Worcesthershure, UK) fitted by a 514 nm laser beam at a fixed angle of 90° and at T=25 °C. Only the samples of LNCs with composition: Lab, 1.028 g; Lip, 0.075 g; Sol, 0.848 g; and hydrodynamic diameter HD=50 nm were used.

By measuring the electro-kinetic properties of LNCs suspensions [6, 7], two populations of LNCs were detected: LNCs containing a small number of lipoid molecules (LNC I) having slightly negative values of the zeta potential (ZP=-5.4 mV) and LNCs with more than a critical content of lipoid molecules (LNC II) having more negative values of the zeta potential (ZP=-20 mV). Both LNCs populations were stable and remained intact in the bulk of their suspensions. A loss of mechanical stability at A/W interface leading to formation of a TG monolayer from the core of the LNCs was observed for the LNC I fraction. The lipid molecules stabilize the LNC's structure. Thus, LNCs with smaller lipoid content are expected to be more unstable at the interface in comparison to LNCs II. From experimental data and structural considerations, the percentage ratio between the two populations in the studied LNCs is evaluated as 85 % (LNCs I) to 15 % (LNCs II) [7].

#### Measurement at the A/W interface

The surface pressure (II) was measured using KSV-2200 (Finland) surface balance, equipped with platinum plate and a Teflon trough with total trough area A=475 cm<sup>2</sup>.

#### Kinetic experiments

Using Exmire microsyringe DPPC monolayers were spread at the water surface from DPPC chloroform solution ( $C_{\text{DPPC}}=1 \text{ mg} \times \text{cm}^{-3}$ ) at low surface pressure  $\Pi < 0.2 \text{ mN} \times \text{m}^{-1}$ . After waiting for about 10 min for solvent evaporation, the monolayers were compressed until one of desired initial surface pressures ( $\Pi=1, 2, 5, 10, 15, 20$ , and 25 mN×m<sup>-1</sup>) was reached. Then, the spreading of LNC dispersion containing LNC I and LNC II from a syringe at level of the water surface on preformed DPPC monolayer was done. In the cases of preformed films from Curosurf or mucus, the monolayers were obtained by spreading Curosurf<sup>®</sup> or mucus from aqueous solutions. After a period of about 30 min needed for the formation of stable surface film, the LNC dispersion was spread.

The following LNCs dispersions with concentrations of the unstable fraction LNC I estimated by means of previously described approach [7] were used:  $C^*=2.5 \times 10^{15}$ ,  $2.5 \times 10^{14}$ , and  $2.5 \times 10^{13}$  (LNC cm<sup>-3</sup>). To follow the disaggregation of unstable LNCs type I, leading to rapid release of TG from the LNCs core and formation of mixed TG/DPPC, TG/Curosurf<sup>®</sup> or TG/mucus monolayers, two kinds of experiments were performed after spreading of LNCs:

- (i) Recording the surface pressure (Π) change with time (t) at constant surface area (A);
- (ii) Recording of the surface area change ( $\Delta A$ ) with time (*t*) at constant surface pressure ( $\Pi$ ).

#### Quasiequilibrium isotherms

The isotherms  $\Pi(A)$  for DPPC, Curosurf<sup>®</sup>, mucus, Labrafac<sup>®</sup> (TG), and mixed TG/DPPC monolayers were also measured and used for comparison and interpretation of kinetic data. After spreading at low surface pressure ( $\Pi$ <0.2 mN m<sup>-1</sup>) the monolayers were compressed by means of barrier movement at sufficiently low speed ( $U_b$ =10.8 cm<sup>2</sup>×min<sup>-1</sup>). The quasiequilibrium isotherms  $\Pi(A)$  for DPPC, Labrafac<sup>®</sup> (TG) and mixed TG/DPPC monolayers obtained during the compression were well reproducible and practically identical at sufficiently low surface pressure with those obtained at decompression. Then, they could be used to obtain the composition TG/DPPC corresponding to the saturation of kinetic data obtained at comparable characteristic times. On the contrary, the dynamic  $\Pi(A)$  isotherms for Curosurf<sup>®</sup> and mucus reveal a net hysteresis (data not shown) and cannot be considered as equilibrium ones.

All measurements were performed at room temperature (T=23 °C).

#### AFM imaging: sample preparation

The monolayers were transferred vertically on mica sheets  $(10 \times 10 \text{ mm})$  by LB deposition where the support was immersed in the subphase before spreading either of DPPC or mucus monolayer. Next, the monolayers were compressed to  $\Pi = 5 \text{ mN} \times \text{m}^{-1}$  and LNC dispersions were deposited. After a 30-min rest, the mica plate was pulled out of the subphase with a transferring rate of 5 mm×min<sup>-1</sup>. The films were placed in exicator for 24 h and before imaging were blown with nitrogen gas for about 1–2 min.

NanoScope Multi Mode V AFM-system (Bruker Inc., Germany) operating in tapping mode at air with silicon cantilevers (Tap300Al-G, Innovative solutions, Bulgaria) with spring constant of  $1.5 \div 15 \text{ N} \times \text{m}^{-1}$  and tip radius about 20 nm. The scan rate was 1 Hz and the images were captured in height and phase modes and only flattened by Nanoscope's software (v.7.30). All samples were typically scanned in several different locations along the mica sheet and representative AFM images were chosen.

#### **Theoretical considerations**

A previously developed theoretical approach allowed us to properly analyze the behavior of LNCs spread at pure A/W interface in case of large spreading quantities and times of order of 60 min [7]. In the present paper, we expand this theoretical framework to the case of LNCs spread on preliminary formed monolayers in order to define both the rate constant  $k_1$ , the maximal degree,  $d_{TG}$ , and the effectiveness of inclusion,  $\varepsilon_{TG}$ , of TG (Labrafac<sup>®</sup>) into mixed DPPC/Labrafac® monolayer at saturation. In addition to that, we establish that the best conditions for comparison the LNCs behavior spread either at pure A/W interface or on top of a preformed monolayer correspond to small spreading LNCs quantities at times less than 1 min. In case of times less than 1 min during first rapid stage of experiment, the main contribution in  $\Pi(t)$ - kinetics is due to rapid disaggregation of unstable LNC I population and instantaneous Labrafac® spreading, which means that all other

processes in general kinetic scheme considered in [7] can be neglected.

The interpretation of experimental kinetic data  $\Delta A(t)_{\Pi=const}$ and  $\Pi(t)_{A=const}$  can be obtained on the basis of following model. A rapid disaggregation of the unstable LNCs I leads to release of confined into capsules TG molecules on the preliminary formed DPPC monolayer with rate constant k<sub>1</sub>. During formation of a mixed TG/DPPC monolayer, the total number n(t) of molecules of both species spread at whole available surface area of the trough is given by following simple expression:

$$n(t) = n_1(t) + n_2(t) = A(t) \left[ \Gamma_1(t) + \Gamma_2(t) \right]$$
(1)

where  $n_1(t)$  and  $n_2(t)$  are the number of TG and DPPC molecules, respectively in the mixed TG/DPPC monolayer formed on the whole available area  $A(t)=A_0 + \Delta A(t)$ ;  $\Gamma_1(t)$  and  $\Gamma_2(t)$ are the surface concentrations of TG and DPPC molecules, respectively in unit number of molecules per unit area.

At constant surface area  $(A_0 = \text{const})$   $n_2^0 \equiv n_2(A_0) = \text{const}$ ;  $\Gamma_2^0 \equiv \Gamma_2(A_0) = \text{const}$  for  $t \ge 0$  during the spreading of TG, Eq. 1 is reduced to:

$$n(t) = n_1(t) + n_2^0 = A_0 \left[ \Gamma_1(t) + \Gamma_2^0 \right]$$
(2)

and

$$\frac{\mathrm{d}n(t)}{\mathrm{d}t} = A_0 \frac{\mathrm{d}\Gamma_1}{\mathrm{d}t}$$

Assuming that the process of disaggregation of unstable capsules type I is rate-determining, while the spreading of TG molecules confined into capsules is instantaneous, the following kinetic equation for the rate of increase of surface concentration of TG molecules in the mixed monolayer can be written as:

$$\frac{\mathrm{d}\Gamma_1}{\mathrm{d}t} = k_1 (1-\theta) m_1^0 N_1 \tag{3}$$

where  $k_1$  is the rate constant of LNCs I disaggregation;  $\theta = \frac{\Gamma_1(t) + \Gamma_2^0}{\Gamma_{\infty}}$  is the degree of surface coverage during formation of the mixed monolayer;  $(1-\theta)$  is the part of free surface area;  $N_I$  is the number of unstable LNCs I per unit area;  $m_1^0$  is the number of TG molecules confined in one unstable LNC of type I;  $m_1^0 N_I$  is the number of TG molecules contained in all LNCs I spread at unit surface.

At the beginning of the process  $\Gamma_1(t)=0$  and  $\theta = \Gamma_2^0 / \Gamma_{\infty}$ ;  $\Gamma_{\infty}$  is surface concentration corresponding to a close

monomolecular packing of DPPC molecules and  $E_2^0 = \frac{d\Pi}{d\Gamma_2^0} \Gamma_2^0$ is the surface elasticity of preliminary formed DPPC monolayer. Then  $\left(\frac{d\Gamma_1}{d\Pi}\right)_{ini} \approx \left(\frac{d\Gamma_2^0}{d\Pi}\right)_{ini} = \frac{\Gamma_2^0}{E_2^0}$  and finally the initial rate of increase of surface concentration of TG in mixed monolayer can be represented as:

$$\left(\frac{\mathrm{d}\Gamma_1}{\mathrm{d}t}\right)_{\mathrm{ini}} = \left(\frac{\mathrm{d}\Gamma_1}{\mathrm{d}\Pi}\right)_{\mathrm{ini}} \left(\frac{\mathrm{d}\Pi}{\mathrm{d}t}\right)_{\mathrm{ini}} = \frac{\Gamma_2^0}{E_2^0} \left(\frac{\mathrm{d}\Pi}{\mathrm{d}t}\right)_{\mathrm{ini}} \tag{4}$$

From Eqs. 3 and 4 one obtains the following expression for the rate constant  $k_1$ :

$$k_1 = \frac{(\mathrm{d}\Pi/\mathrm{d}t)_{\mathrm{ini}} \, \Gamma_2^0}{E_2^0 (1-\theta) m_1^0 N_{_J}} \tag{5}$$

A maximal degree of inclusion of TG,  $d_{TG}$ , into mixed monolayer at saturation can be defined as the ratio between number of TG and DPPC species in the mixed monolayer:

$$d_{\rm TG} = \frac{n_1^{\rm sat}}{n_2^0} = \frac{\Gamma_1^{\rm sat}}{\Gamma_2^0} \tag{6}$$

The corresponding effectiveness of inclusion,  $\varepsilon_{TG}$ , is given as the ratio between the numbers of TG molecules included in the surface film at saturation as a result of LNC I disaggregation and those contained in all spread LNCs:

$$\varepsilon_{\rm TG} = \frac{n_1^{\rm sat}}{m_1^0 N_I A} = \frac{\Gamma_1^{\rm sat}}{m_1^0 N_I} \tag{7}$$

When all TG molecules contained in all spread LNCs are included in the surface monolayer, then  $\varepsilon_{TG}=1$ .

The values of  $\Gamma_1^{\text{sat}}$  and then of  $d_{\text{TG}}$  and  $\varepsilon_{\text{TG}}$  can be estimated in two ways. The first one considers the change of surface pressure from  $\Pi_0$  DPPC to  $\Pi_{\text{sat TG/DPPC}}$  which is then compared with independently measured isotherm network for obtaining the composition of mixed TG/DPPC monolayer,  $x_{\text{TG}}$ , resp.  $\Gamma_1^{\text{sat}}$  corresponding to the saturation.

In the second one, after integrating Eq. 4 from t=0 to  $t_{sat}$  and using the experimental values for  $E_2^0$ , we can also estimate the values for  $\Gamma_1^{sat}$  as follows:

$$\Gamma_1^{\text{sat}} \cong \int_0^{t_{\text{sat}}} \frac{\Gamma_2^0}{E_2^0} \left(\frac{\mathrm{d}\Pi}{\mathrm{d}t}\right)_{\text{ini}} \mathrm{d}t \tag{8}$$

In the case when LNCs spread on pure A/W interface form a pure TG monolayer ( $\Gamma_2=0$ ),  $\Gamma_1^{sat}$  can be determined by comparison with the isotherm  $\Pi(A=1/\Gamma)$  for pure TG monolayer

а

П (mN.m<sup>-1</sup>)

b

П (mN.m<sup>-1</sup>)

60

50

40

The rate constant  $k_1$  in this case of formation of pure TG monolayer can be estimated from Eq. 3 as follows:

$$k_1 = \frac{\left(\Gamma_{\rm TG}^{\rm sat}/t_{\rm sat}\right)}{\left(1 - \theta_{\rm TG}^{\rm sat}\right) m_1^0 N_1} \tag{9}$$

At constant surface pressure ( $\Pi$ =const)

If assuming an ideal behavior of the formed mixed TG/ DPPC monolayer, we can consider the initial surface area  $A_0$ as covered by DPPC molecules and the observed surface area change ( $\Delta A$ ) covered only by TG molecules with constant surface densities  $\Gamma_1^0$  and  $\Gamma_2^0$ , respectively. Then, the Eq. 1 is reduced to:

$$n(t) = n_1(t) + n_2^0 = \Delta A(t) \Gamma_1^0 + A_0 \Gamma_2^0$$
(10)

From Eq. 10, we can determine also the values for maximal degree of inclusion  $d_{TG}$  of TG in DPPC monolayer as follows:

$$d_{\rm TG} = \frac{\Delta A_{\rm sat} \Gamma_1^0}{A_0 \Gamma_2^0} \tag{11}$$

and the corresponding effectiveness of inclusion  $\varepsilon_{TG}$  as:

$$\varepsilon_{\rm TG} = \frac{\Delta A_{\rm sat} \, \Gamma_1^0}{m_1^0 N_I A} \tag{12}$$

Finally, we have to underline that the theoretical analysis of experimental data assumes that when an aqueous dispersion of LNC has been spread on the top of existing monolayer of either DPPC, Curosurf<sup>®</sup>, or mucus the spreading procedure is quantitative, i.e., there is no loss of material in bulk phase during the spreading process.

#### **Results and discussion**

Isotherms of the studied monolayers

The quasiequilibrium isotherms  $\Pi(A=1/\Gamma)$  for various compositions  $X_{Lab}$  of model mixed Lab/DPPC monolayers were presented in Fig. 1a. They were used to obtain the values for  $\Gamma_{\text{Lab}}$  at saturation after comparison with the  $\Pi(t)$  kinetic data.

The Labrafac® is a mixture of C:8 and C:10 TG and curve 7 at Fig. 1a corresponds to  $\Pi(A)$  isotherm of pure Labrafac<sup>®</sup> monolayer in liquid-expanded (LE) state with collapse at  $\Pi \approx 22 \text{ mN} \times \text{m}^{-1}$  a value which is in the same order of

30 20 10 0 20 40 60 80 100 120 140 A ( $Å^2$ .molecule<sup>-1</sup>) 60 50 40 30 20 10 0 2000 4000 6000 0 8000 A  $(cm^2.mg^{-1})$ 

0.2 0.4 0.6 0.8 1311

**Fig. 1** a Surface pressure  $(\Pi)$ /mean area  $(\overline{A})$  per molecule isotherms of mixed Labrafac®/DPPC monolayers: curve 1, pure DPPC; curve 7, pure Labrafac®; curves 2-6, mixed Lab/DPPC monolayers with different molecular ratio X<sub>Lab</sub>: 2, X<sub>Lab</sub>=0.11; 3, X<sub>Lab</sub>=0.15; 4, X<sub>Lab</sub>=0.25; 5,  $X_{\text{Lab}}=0.50; 6, X_{\text{Lab}}=0.75 \text{ (inset, } \overline{A} \text{ versus } X_{\text{Lab}} \text{ for a mixed Labrafac}^{\mathbb{R}}/$ DPPC monolayers at  $\Pi$ =2 mN×m<sup>-1</sup> (black square); 5 mN×m<sup>-1</sup> (blue square); 10 mN×m<sup>-1</sup> (red square); 15 mN×m<sup>-1</sup> (pink square); **b**  $\Pi$ (A) isotherms of DPPC, Curosurf<sup>®</sup>, and mucus expressed in  $cm^2 \times mg^{-1}$ : curve 1, DPPC; curve 2, Curosurf®; and curve 3, mucus)

magnitude as the collapse surface pressure of both TG components of Labrafac®. Curve 1 in Fig. 1a corresponding to  $\Pi(A)$  isotherm of pure a DPPC monolayer demonstrates the well-known liquid-expanded/liquid-condensed (LE/LC) phase transition at  $\Pi = 6 \text{ mN} \times \text{m}^{-1}$ . The compositions with increasing DPPC quantities (starting at Fig. 1a from curve 6 with  $X_{\text{DPPC}}=0.25$  to curve 2 with  $X_{\text{DPPC}}=0.89$ ) show that LE/ LC phase transition was influenced by the presence of Labrafac<sup>®</sup>. The inset in Fig. 1a representing the mean molecular area  $\overline{A}$  as a function of molar ratio  $X_{\text{Lab}}$  indicates miscibility with a slightly practically negligible positive deviation from additivity.

The peculiar observation in Fig. 1a concerning the shift of curves 2, 3, and 4 to molecular areas smaller than that of pure

DPPC does not really make sense. However, this result can be explained with the presence of soluble components in Labrafac<sup>®</sup>, which are squeezed out from the monolayer at larger surface pressures. Such effect is not taken into account in the calculation of mean area  $\overline{A}$  of mixed Lab/DPPC monolayers.

The  $\Pi(A)$  isotherms of the three studied preformed monolayers—DPPC, Curosurf<sup>®</sup>, and mucus, are also presented in Fig. 1b. Considering the complex composition of Curosurf<sup>®</sup> and mucus, the areas in Fig. 1b were expressed in square centimeters per milligram. For comparison, the isotherm of DPPC monolayer in Fig. 1b was also transposed in the same units.

We have to note that in case of monolayer formed from natural mixture such as mucus fluid, obviously it is difficult to accurately determine whether this material will retain its structure and properties after spreading at A/W interface. Nevertheless, it may be assumed that the network formed by negatively charged mucin fibers could preserve and maintain an interfacial organization not much different from that, which exists in the mucus layer covering luminal surface of gastrointestinal tract.

#### Spreading of different LNCs quantities at pure A/W interface

The  $\Pi(t)_{A=\text{const}}$  results obtained after spreading of different LNCs quantities  $N_I$  (column 1 of Table 1) on pure A/W interface are presented in Fig. 2. The observed large jump between curves 3 and 4 at Fig. 2 is due to the large difference of surface compressibility modulus  $C_s^{-1} \equiv E$  $(E_{\text{Lab}(3)}=35 \text{ mN} \times \text{m}^{-1}; E_{\text{Lab}(4)}=71 \text{ mN} \times \text{m}^{-1})$ . We have to note that the surface pressure at saturation ( $\Pi_{\text{sat}}$ ) corresponding to close packing of Labrafac<sup>®</sup> monolayer cannot be reached within the experimental framework of spreading LNC quantities  $N_I$  (i.e., from  $0.3 \times 10^{10}$  to  $3.2 \times 10^{10}$ LNC cm<sup>-2</sup>). For comparison, in the previous paper [7], the maximal saturation in Labrafac<sup>®</sup> monolayer at  $\Pi=$ 22 mN×m<sup>-1</sup> corresponded to collapse surface pressure, which has been obtained for  $N_I=1.2 \times 10^{11} \text{ LNC} \times \text{cm}^{-2}$ .



**Fig. 2** Evolution of surface pressure (II) with time (*t*) at A=constant after spreading different LNCs quantities  $N_{\rm I}$  (in LNC×cm<sup>-2</sup>) on pure A/W interface: *curve 1*,  $N_I$ =0.3.10<sup>10</sup>; *curve 2*,  $N_I$ =0.5×10<sup>10</sup>; *curve 3*,  $N_I$ =0.8×10<sup>10</sup>; *curve 4*,  $N_I$ =1.6×10<sup>10</sup>; *curve 5*,  $N_I$ =3.2×10<sup>10</sup>

The slight delay of surface pressure increase in curves 1 and 3 and the slight difference in the course of kinetic curve 4 (Fig. 2) recorded in the first 10 s has no physical meaning because this time is comparable with the time of LNCs spreading procedure.

By comparing the values of surface pressure at saturation  $\Pi_{sat}$  (column 3, Table 1) with  $\Pi(A)$  isotherm for Labrafac<sup>®</sup> monolayer (Fig. 1a, curve 7), the values for  $\Gamma_{TG}^{sat}$  (column 4, Table 1) and according to Eq. 7 determined the values of effectiveness of inclusion  $\varepsilon_{TG}$  (column 7, Table 1). Using the experimental data for  $\Gamma_{TG}^{sat}/t_{sat}$ , for  $\theta_{TG}^{sat} = \Gamma_{TG}^{sat}/\Gamma_{\infty}$  from  $\Pi(A)$  isotherm of Labrafac<sup>®</sup> and Eq. 9 applied for the case of LNCs spreading on pure A/W interface, the values of the rate constant  $k_1$  of LNCs disaggregation (column 8, Table 1) are also estimated.

At smallest LNCs spreading quantity  $N_I = 0.3 \times 10^{10}$  LNC×cm<sup>-2</sup>, the monolayer at saturation is formed from all TG molecules contained in all spread LNCs type I and the effectiveness of disaggregation of the LNCs is 100 % ( $\varepsilon_{\rm TG}$ =1). When the LNCs spreading quantity  $N_I$  increases, then the effectiveness,  $\varepsilon_{\rm TG}$ , at the saturation

Table I Spreading of different	
LNCs quantities on pure A/W	
interface	

$\frac{1}{\left(\frac{\text{LNC}}{\text{cm}^2}\right)}$	$2 \\ \left(\frac{m_1^0 N_{\rm I}}{{\rm cm}^2}\right)$	$3 \\ \left(\frac{mN}{m}\right)$	$\begin{array}{c} 4 \\ \Gamma_{TG}^{sat} \\ \left( \frac{molTG}{cm^2} \right) \end{array}$	5 $\theta_{\rm TG}^{\rm sat}$	$\begin{array}{c} 6 \\ \Delta \Gamma_{TG} / \Delta t \\ \left( \frac{molTG}{cm^2 \ min} \right) \end{array}$	$\frac{7}{\varepsilon_{\mathrm{TG}}} = \frac{\Gamma_{\mathrm{TG}}^{\mathrm{sat}}}{m_1^0 N_I}$	$egin{array}{c} k_1 \ (\min^{-1}) \end{array}$
$0.3 \times 10^{10}$	$0.69 \times 10^{14}$	0.93	$0.68 \times 10^{14}$	0.60	$1.5 \times 10^{14}$	1.00	5.8
$0.5 \times 10^{10}$	$1.15 \times 10^{14}$	2.20	$0.77 \times 10^{14}$	0.70	$2.2 \times 10^{14}$	0.67	6.8
$0.8 \times 10^{10}$	$1.85 \times 10^{14}$	4.60	$0.83 \times 10^{14}$	0.83	$2.1 \times 10^{14}$	0.45	6.7
$1.6 \times 10^{10}$	$3.68 \times 10^{14}$	14.0	$1.00 \times 10^{14}$	0.91	$5.0 \times 10^{14}$	0.27	15.1
$3.2 \times 10^{10}$	$7.36 \times 10^{14}$	15.9	$1.04 \times 10^{14}$	0.95	$5.2 \times 10^{14}$	0.14	14.1

decreases. This effect could be related either to partial stabilization of the LNCs given their interaction with the already formed TG monolayer at the latter stage of spreading or/and to the fact that part of the free surface  $(1-\theta)$  diminishes from 0.40 to 0.05.

The estimated values of the rate of LNCs disaggregation  $k_1$  are ranging between 5 and 15 min<sup>-1</sup>.

To understand better the loss of LNCs mechanical stability at pure A/W interface or covered by preformed monolayer models of membranes, the small spreading LNCs quantities showing  $\varepsilon_{TG}$ =1 at pure A/W interface were studied and compared. One of the reasons to study the behavior of LNCs spread at small quantities was the observation that at smallest LNCs spreading quantities, the monolayer at saturation is formed from all Labrafac<sup>®</sup> molecules contained in all spread LNCs type I, i.e., there is no loss of material during the spreading procedure of LNC at A/W interface.

Spreading of small LNCs quantities on preliminary formed DPPC monolayer

The  $\Pi$  (*t*)<sub>*A*=const</sub> results obtained after spreading of small LNCs quantities  $N_I$  from 0.41 to  $0.84 \times 10^{10}$  and from 0.05 to  $0.08 \times 10^{10}$  LNC cm<sup>-2</sup>, corresponding to ratio 1:1, respectively 1:10 between the TG molecules confined in all LNCs and DPPC molecules in the previously spread DPPC monolayer at various surface pressures  $\Pi = 1, 2, 5, 10, 15, 20$ , and  $25 \text{ mN} \times \text{m}^{-1}$  are presented in Fig. 3a. The observed large jump between curves 3 and 4 seems to take place at the same surface pressure where DPPC undergoes a phase transition. The heterogeneous structure of monolayer probably plays a decisive role in interfacial mechanical stability of LNCs as it would be seen and discussed later from AFM data, where LNCs are attached preferably at the boundary between the two coexisting DPPC phases. In addition, as observed in Fig. 2, the jump may also be related to differences of the surface compressibility modulus  $C_s^{-1} \equiv E (E_{\text{DPPC}(3)} = 39 \text{ mN m}^{-1};$  $E_{\text{DPPC}(4)} = 68 \text{ mN m}^{-1}$ ).

The values for  $I_{TG}^{sat}$  are obtained either by comparison with the mixed TG/DPPC monolayer  $\Pi(A)$  isotherms presented in Fig. 1 or from experimental values for the surface elasticity  $E_{DPPC}$  and Eq. 8. The values for  $d_{TG}$ ,  $\varepsilon_{TG}$ , and  $k_1$  corresponding to the same spreading LNCs quantities obtained from both procedures are practically identical. The summarized results for  $d_{TG}$ ,  $\varepsilon_{TG}$ , and  $k_1$  obtained from experimental data  $\Pi(t)_{A=const}$  using both approaches are presented in Table 2.

The  $\Delta A(t)_{\Pi=const}$  results obtained after spreading of LNCs quantities  $N_I=0.55\div0.84\times10^{10}$  and 0.05 to  $0.08\times10^{10}$  LNC cm<sup>-2</sup> are presented at Fig. 3b. The interpretation of  $\Delta A(t)$  at  $\Pi_0=$ const is based on Eqs. 10, 11, and 12. The obtained values for  $d_{\rm TG}$  and  $\varepsilon_{\rm TG}$  differ slightly from those obtained by two previous ways of calculation at  $A_0=$ const



**Fig. 3** a  $\Pi(t)_{A=const}$  after spreading of LNCs on preliminary formed DPPC monolayer; molecular ratio Lab/DPPC (1:1) at various surface pressure: *curve 1*,  $\Pi=1 \text{ mN} \times \text{m}^{-1}$ ; *curve 2*,  $\Pi=2 \text{ mN} \times \text{m}^{-1}$ ; *curve 3*,  $\Pi=5 \text{ mN} \times \text{m}^{-1}$ ; *curve 4*,  $\Pi=10 \text{ mN} \times \text{m}^{-1}$ ; *curve 5*,  $\Pi=15 \text{ mN} \times \text{m}^{-1}$ ; *curve 6*,  $\pi=20 \text{ mN} \times \text{m}^{-1}$ ; and *curve 7*,  $\Pi=25 \text{ mN} \times \text{m}^{-1}$  (*inset* corresponds to molecular ratio Lab/DPPC 1:10 at surface pressures  $\Pi=5$ , 10, 15, 20, and  $25 \text{ mN} \times \text{m}^{-1}$ ); **b**  $\Delta A(t)_{II=const}$  after spreading of LNCs on preliminary formed DPPC monolayer; molecular ratio Lab/DPPC (1:1) at various surface pressure: *curve 1*,  $\Pi=5 \text{ mN} \times \text{m}^{-1}$ ; *curve 2*,  $\Pi=10 \text{ mN} \times \text{m}^{-1}$ ; *curve 3*,  $\Pi=15 \text{ mN} \times \text{m}^{-1}$ ; *curve 4*  $\Pi=20 \text{ mN} \times \text{m}^{-1}$ ; and *curve 5*,  $\Pi=25 \text{ mN} \times \text{m}^{-1}$  (*inset* corresponds of molecular ratio Lab/DPPC 1:10 at the same surface pressures)

(data not shown). Nevertheless, the same dependence on the LNCs spread quantities is observed. As far as the observed small differences are concerned it is important to note that the process of LNCs destabilization and disaggregation occurs in better defined thermodynamical conditions at  $\Pi_0$ =const than at  $A_0$ =const.

**Fig. 4**  $\Delta A(t)_{II=\text{const}}$  (**a**) and  $\Pi(t)_{A=\text{const}}$  (**c**) after spreading of LNCs on preformed Curosurf<sup>®</sup> monolayer at various surface pressures: *curve 1*,  $\Pi=5 \text{ mN} \times \text{m}^{-1}$ ; *curve 2*,  $\Pi=10 \text{ mN} \times \text{m}^{-1}$ ; *curve 3*,  $\Pi=15 \text{ mN} \times \text{m}^{-1}$ ; *curve 4*,  $\Pi=20 \text{ mN} \times \text{m}^{-1}$ ; *and curve 5*,  $\Pi=25 \text{ mN} \times \text{m}^{-1}$ ;  $\Delta A(t)_{II=\text{const}}$  (**b**) and  $\Pi(t)_{A=\text{const}}$  (**d**) after spreading of LNCs on preformed mucus monolayer at the same surface pressures (\Pi)



Comparison between the obtained results for LNCs spread at pure A/W interface and on a DPPC monolayer

The comparison between the obtained results for LNCs spread at pure A/W interface (Table 1) and on the top of preliminary formed DPPC monolayer (Table 2) suggests that the interaction between LNCs and DPPC monolayer is responsible for a partial LNCs stabilization. For example, at approximately the same  $N_I$ =0.3÷0.41×10<sup>10</sup> LNC×cm<sup>-2</sup>,  $\Pi$ =0.93÷1.0 mN× m<sup>-1</sup>, and  $\theta$ =0.50÷0.60, the effectiveness of LNC spreading  $\varepsilon_{\rm TG}$  given as the ratio between the numbers of TG molecules included in the surface film at saturation as a result of LNC I disaggregation and those contained in all spread LNCs and rate constant of LNC disaggregation  $k_1$  are  $\varepsilon_{\rm TG}$ =1 and  $k_1$ = 5.8 min<sup>-1</sup> on pure A/W interface while at preformed DPPC monolayer the effectiveness and the rate of disaggregation

1	2	2	4	5	6
1	2	5	+	5	0
$\Pi^0_{\text{DPPC}}$	(LNC)	$\theta_{\text{DPPC}} = \frac{\Gamma^0}{\Gamma_\infty}$	$d_{\rm TG} = \frac{\Gamma_{\rm TG}^{\rm sat}}{\Gamma_{\rm DPPC}^0}$	$\varepsilon_{\mathrm{TG}} = \frac{\Gamma_{\mathrm{TG}}^{\mathrm{sat}}}{m_1^0 N_I}$	$k_1$ $(min^{-1})$
$\left(\frac{\mathrm{IIIN}}{\mathrm{m}}\right)$	$\left(\frac{\mathrm{d}\mathrm{i}\mathrm{v}\mathrm{e}}{\mathrm{c}\mathrm{m}^2}\right)$		bite	Ĩ	(11111 )
1	$0.41 \times 10^{10}$	0.50	0.510	0.51	2.7
2	$0.44 \times 10^{10}$	0.52	0.430	0.43	2.9
5	$0.05 \times 10^{10}$	0.65	0.030	0.32	1.9
	$0.55 \times 10^{10}$	0.65	0.280	0.27	1.8
10	$0.08 \times 10^{10}$	0.78	0.020	0.18	1.2
	$0.76 \times 10^{10}$	0.78	0.230	0.17	2.3
15	$0.08 \times 10^{10}$	0.83	0.010	0.11	1.1
	$0.76 \times 10^{10}$	0.83	0.120	0.12	2.4
20	$0.08 \times 10^{10}$	0,88	0.005	0.05	0.9
	$0.84 \times 10^{10}$	0.88	0.090	0.09	1.9
25	$0.08 \times 10^{10}$	0.91	0.001	0.01	0.2
	$0.75 \times 10^{10}$	0.91	0.030	0.03	1.0

**Table 2** Spreading of smallLNCs quantities at  $A_0$ =const onpreformed DPPC monolayers

are smaller  $\varepsilon_{TG}=0.51$  and  $k_1=2.7 \text{ min}^{-1}$ . The same effect—larger  $\varepsilon_{TG}$  and  $k_1$  for pure A/W interface than for DPPC monolayer is observed also for studied quantities  $N_I=0.5\div0.55\times10^{10}$  LNC×cm<sup>-2</sup> ( $\varepsilon_{TG}=0.67$  and  $k_1=6.8 \text{ min}^{-1}$  versus  $\varepsilon_{TG}=0.27$  and  $k_1=1.8 \text{ min}^{-1}$ ) and  $N_I=$  $0.8\div0.76\times10^{10}$  LNC×cm<sup>-2</sup> ( $\varepsilon_{TG}=0.45$  and  $k_1=6.7 \text{ min}^{-1}$ versus  $\varepsilon_{TG}=0.17$  and  $k_1=2.3 \text{ min}^{-1}$ ). Thus, the comparison suggests that the observed difference in the effectiveness of LNC spreading,  $\varepsilon_{TG}$ , and rate of disaggregation,  $k_1$ , for LNCs spread at the pure A/W interface and on preformed DPPC monolayer is due to the LNC partial stabilization and not to the lack of free surface area  $(1-\theta)$ .

As it was shown in "Spreading of different LNCs quantities at pure air-water interface" when LNCs spreading quantity,  $N_{I}$ , increase the observed decrease of the effectiveness,  $\varepsilon_{TG}$ , could be in principle related either to the LNCs partial stabilization or to the lack of free surface area  $(1-\theta)$ . The idea that a partial LNCs stabilization occurs as a result of the interaction with the A/W surface covered by DPPC molecules is also confirmed by the fact that at the same part of free surface area  $(1-\theta)$ , the same effectiveness  $\varepsilon_{TG}$  and rate of disaggregation  $k_1$  are obtained with 10-fold larger LNCs spread quantities  $N_I$ and degree of inclusion  $d_{TG}$ . For example, at the same  $\Pi = 5 \text{ mN} \times \text{m}^{-1}$  (Table 2) and the same part of free surface area  $(1-\theta)=0.35$  for  $N_I=0.05\times 10^{10}$  LNC×cm<sup>-2</sup> and for  $N_I=0.05\times 10^{10}$  LNC×cm<sup>-2</sup>  $0.55 \times 10^{10}$  LNC×cm<sup>-2</sup>, the obtained values  $d_{\rm TG}$ =0.030,  $\varepsilon_{\rm TG}$ = 0.32,  $k_1 = 1.9 \text{ min}^{-1}$  and  $d_{TG} = 0.280$ ,  $\varepsilon_{TG} = 0.27$ ,  $k_1 = 1.8 \text{ min}^{-1}$ respectively are practically the same.

Spreading of small LNCs quantities on preliminary formed Curosurf<sup> $\mathbb{R}$ </sup> and mucus monolayers

The  $\Delta A(t)_{\Pi=\text{const}}$  and  $\Pi(t)_{A=\text{const}}$  results obtained after spreading of 0.028 mg TG (1.43×10<sup>12</sup> LNC) on the whole monolayer area covered with 0.071 mg Curosurf<sup>®</sup> are presented in Fig. 4a, c. The interpretation of experimental data is as described above and the obtained results for  $d_{\text{TG}}$ ,  $\varepsilon_{\text{TG}}$ , and  $k_1$  are presented in Table 3. The ratio between TG and Curosurf<sup>®</sup> is expressed as mg per mg because of the complex composition of Curosurf<sup>®</sup>. To compare  $d_{\text{TG}}$  and  $\varepsilon_{\text{TG}}$ , the units for DPPC monolayer were also transposed in mg instead of molecules. The obtained values for effectiveness  $\varepsilon_{\text{TG}}$ , of LNCs disaggregation on Curosurf<sup>®</sup> monolayer were close to those on DPPC monolayer. This finding seems logical because DPPC is the main component of Curosurf<sup>®</sup>.

The results  $\Delta A(t)_{\Pi=const}$  and  $\Pi(t)_{A=const}$  obtained after spreading of 0.011 mg TG ( $1.43 \times 10^{12}$  LNC) on the whole monolayer area covered with 0.17 mg mucus are presented in Fig. 4b, d. The values for  $d_{TG}$ ,  $\varepsilon_{TG}$ , and  $k_1$  obtained by means of the above-described approach are summarized and presented also in Table 3.

At the same lateral surface pressure, the effectiveness  $\varepsilon_{TG}$  of LNCs disaggregation on mucus monolayer decreases and mechanical stability increases in comparison with those on DPPC monolayer.

As far as the molecular nature of LNCs-model membrane monolayers interaction is concerned, the obtained results



Fig. 5 Determination of MIP of TG from disaggregated LNCs into DPPC (a), Curosurf® (b), and mucus (c) monolayers

about LNCs mechanical stability viewed at phenomenological level by  $\varepsilon_{TG}$  seem to be in accordance with the common understanding about the role of electrostatic "particle-monolayer" interactions [24]. The interaction between LNCs and monolayer model membrane determining the mechanical stability (or instability) is a result of superposition of various electrostatic and hydrophobic interactions. The observation that the LNCs are more stable when spread at mucus rather than on top of DPPC or Curosurf® monolayer could be explained by the electrostatic interaction between them. Actually, the particle shell of LNCs is slightly negatively charged. From measurements of electro-kinetic properties of LNC suspensions [6, 7], the values for zeta potential ZP=-5.4 mV for LNC I fraction without or containing only small number of lipoid molecules and respectively ZP=-20 mV for the LNC II fraction with more than a critical content of lipoid molecules were obtained. The LNCs have negative zeta potential [25] very close to the ZP of other carriers containing PEG such as liposomes because the LNCs shells contain hydrophilic surfactants with PEG chains [26]. The phosphate groups of phospholipid molecules presented in larger quantity in LNC II fraction should increase the electro-negativity if we assume that they can deliver these charges at the external part of shell. On the other side, the

**Table 3** Spreading of small LNCs quantities at  $\Pi_0$ =const on preformed DPPC, Curosurf<sup>®</sup>, and Mucus monolayers: Spread on the whole surface area quantities: DPPC, 0.04 mg; TG, 0.028 mg ( $m_1^0$ NA<sub>0</sub> = 1.43×10<sup>12</sup> LNC); Curosurf<sup>®</sup>, 0.071 mg; TG, 0.011 mg ( $m_1^0$ NA<sub>0</sub>=0.6.10<sup>12</sup> LNC); mucus, 0.17 mg; TG, 0.011 mg;  $I_2^0$  is the surface density at initial surface area  $A_0$  for DPPC, Curosurf<sup>®</sup> and mucus, respectively

$\Pi^0 (\mathrm{mN \ m}^{-1})$	$d_{ m TG}=rac{arGamma_{ m TG}^{ m sat}}{arGamma_2^0}$	$arepsilon_{\mathrm{TG}}=rac{I_{\mathrm{TG}}^{\mathrm{sat}}}{m_{1}^{0}N_{I}}$	$k_1 (\min^{-1})$
DPPC			
5	0.28	0.40	1.8
10	0.19	0.26	2.3
15	0.12	0.18	2.4
20	0.08	0.11	1.9
25	0.04	0.05	1.0
Curosurf®			
5	0.20	0.49	6.8
10	0.13	0.33	7.1
15	0.08	0.21	4.6
20	0.06	0.15	3.5
25	0.03	0.07	1.9
Mucus			
5	0.03	0.19	6.9
10	0.02	0.14	4.9
15	0.01	0.07	3.8
20	0.01	0.04	5.5
25	0.00	0.01	0.0

mucus is composed of negatively charged mucin fibers [27]. These negative charges in the plane of the monolayer could be partly responsible for the mucin-LNC electrostatic repulsion leading to LNCs stabilization.

In opposition to all that, the monolayers of DPPC as well as Curosurf<sup>®</sup>, where the main component is DPPC, are composed of neutral molecules. Thus, no electrostatic repulsion between LNCs and either of DPPC or Curosurf<sup>®</sup> monolayers are expected.

Maximal insertion pressure of the studied systems

The degree of inclusion strongly depends on surface pressure of the preliminary spread monolayer and it is usually characterized by the so-called maximal insertion pressure (MIP) [28].

From the experimental  $\Pi(t)$  data presented in Figs. 3a and Fig. 4c, d represented in the scale increase of surface pressure  $\Delta\Pi$  as function of the initial surface pressure  $\Pi_i$  Fig. 5 the following values for MIP: MIP=31.1 and 31.7 mN×m<sup>-1</sup> were obtained for DPPC and Curosurf<sup>®</sup> monolayers, respectively, and MIP=23.6 mN×m<sup>-1</sup> for mucus monolayer. The larger values of MIP imply that the penetration of TG into DPPC and Curosurf<sup>®</sup> monolayer is easier than into the mucus monolayer.

#### AFM imaging

It is well-established that DPPC monolayer undergoes a firstorder phase transition from liquid-expanded (LE) phase to liquid-condensed (LC) phase, as a result of process driven by a competition between line tension and electrostatic forces [29]. The LE/LC transition in monolayers has been structurally well characterized on micrometer scale level by fluorescence and Brewster angle microscopies [30]. It was also shown that the coexisting lateral structure of phospholipid monolayers can be preserved after their transfer via LB technique on solid mica supports for AFM imaging [31]. Following the same protocol as in [31], we imaged the state of preformed DPPC and mucus monolayers in presence of LNCs.

At Fig. 6a are presented AFM image showing the topography of DPPC/LNCs film with embedded into DPPC monolayer two LNCs. Scanning over different areas of the films showed that only few LNCs entrenched in DPPC monolayer were identified. The LE/LC phase separation within DPPC monolayer is confirmed by high resolution topography image shown as an inset in Fig. 6a where the darker areas represent LE phase and brighter LC phase. It follows from the measured height difference between the coexisting phases which is  $\Delta h \approx 0.6$  nm. It is in agreement with the length difference between DPPC carbon chains when the monolayer is in LE and LC phases. Figure 6a also demonstrates that two-phase lipid organization favors the attachment of LNCs at the boundary between LE and LC phases. A similar preference of lipolytic enzyme PhospholipaseA<sub>2</sub> (PLA<sub>2</sub>) to accumulate along the LC/LE domain boundaries is reported [32]. The measured height of LNCs about  $18\pm2.0$  nm indicates for partially disintegration of LNCs which initial diameter before spreading on A/W interface was about 50 nm.

At Fig. 6b are presented AFM images and cross-section graphs of pure mucus monolayer transferred on mica support at 5 mN×m<sup>-1</sup> (left image) and preformed mucus monolayer with LNCs (right image) where the differences in their



**Fig. 6** a AFM topography images of LB films transferred from DPPC monolayer compressed to  $\Pi = 5 \text{ mN} \times \text{m}^{-1}$  followed by deposition of LNC dispersion (*inset* is area zoom with z=10 nm). Beneath, cross-section of LNCs, embedded into DPPC; **b** AFM topography images of LB films of pure mucus monolayer (*left image*) and AFM image of LB film from preformed mucus monolayer with LNCs (*right image*). The *insets* are particular cross-sections

morphologies are easily distinguishable. The morphology of pure mucus films show only an existence of embedded proteins as a part of the mucus content, which have sizes of about 5 nm, a typical value for globular proteins or enzymes [33, 34]. The AFM image of mixed mucus/LNCs LB films shows an abundance of 3D-structures with sizes of about  $3\div5$  nm which are grafted in mucus layer. There is also a population of 3D-structures which could be LNCs embedded in the mucus layer because of their much bigger size ( $10\div15$  nm) close to the size of LNCs imaged in DPPC/LNCs system. The differences in the morphologies between pure mucus and mucus/ LNCs films are also confirmed from the roughness analysis which shows the estimated mean roughness of 0.179 nm for mucus and about 0.478 nm for mucus/LNCs.

#### Conclusions

A loss of mechanical stability of LNCs spread at pure A/W interface or at preformed model membrane monolayers is observed. The LNCs undergo a partial destabilization, disaggregation, and formation of pure or mixed surface films. The comparison between the effectiveness of spreading ( $\varepsilon_{TG}$ ) and the rate of LNC disaggregation  $(k_1)$  at pure A/W interface and at preformed DPPC monolayer indicates a partial LNCs stabilization due to LNCs-DPPC interaction. The comparison between effectiveness of inclusion of TG ( $\varepsilon_{TG}$ ) into model membrane monolayers of DPPC, Curosurf®, and mucus shows that the LNCs spread on mucus surface layer are more stable and interact stronger with the monolayer than if they are spread on DPPC and Curosurf® monolayers. These findings are in accordance with AFM images showing more 3D structures grafted in mucus layer than in DPPC monolayer. As far as the molecular nature of LNCs-model membrane monolayer interactions are concerned, additional approaches and techniques are needed to develop a conclusive molecular model of the behavior of LNCs spread on various model membrane systems.

From physiological point of view, we hope that such approach could provide some new insights and suggestions about the behavior of LNCs at membranes interfaces. The choice of Curosurf<sup>®</sup> and mucus monolayers as membrane models was mainly provoked by the attempts which are now in progress various liposoluble drugs to be administrated via LNCs towards alveolar surface and luminal surface of gastro-intestinal tract thus proving their effectiveness as drug delivery systems. Obviously, it is a difficult task to bridge the existing gap between monolayer model system and LNCs behavior in physiological medium. Although, in the future we intend to find a way to correlate the results obtained from the developed 2D model of LNCs–membrane interaction with the results from 3D drug delivery model systems.

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